

A sharp interface method for two-phase compressible flows at low-Mach regime

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Journée Des Doctorants du LIMSI



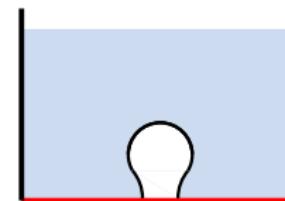
Introduction

Flows of interest

- Compressible gas & quasi-incompressible liquid
- surface tension & heat transfer & phase change

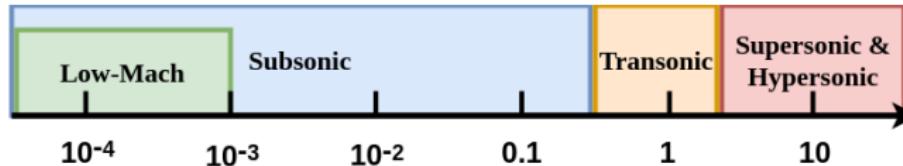


ocean wave



bubble growth near a heated plate

- liquid: low-Mach regime



Mach number ($Ma=u/c$) & Flow Regimes

Introduction

Previous studies

- incompressible liquid model [Tanguy et al., 2014] [Sato and Ničeno, 2013]
 - robust, efficient
 - **without compressible effects**
- compressible liquid model [Fechter et al., 2017][Faccanoni et al., 2012]
 - take compressible effects into account
 - **excessive numerical diffusion** Error $\propto 1/Ma$
 - **time step limitation** $\Delta t = \Delta x / \max(c + u)$

Introduction

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Objective

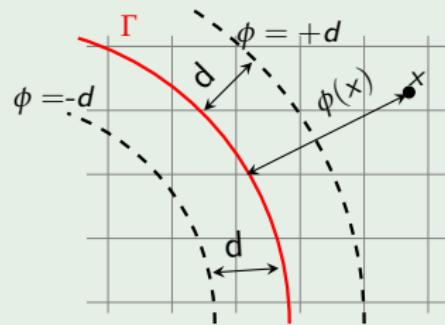
develop a compressible solver for both liquid and gas:

- Accurate at low-Mach regime
- Uniform time step with respect to the Mach number
- Well-balanced discretization of sources terms (gravity, surface tension)
- An accurate high-order method to capture the interface

Interface description

Level-Set method [Osher and Sethian, 1988]

- ϕ : signed distance function
- $|\nabla\phi| = 1$
- Normal vector $\mathbf{n} = \frac{\nabla\phi}{|\nabla\phi|} \Big|_{\phi=0}$
- interface Curvature $\kappa = \nabla \cdot (\frac{\nabla\phi}{|\nabla\phi|})$
- Level-Set advection $\frac{\partial\phi}{\partial t} + \mathbf{u} \cdot \nabla\phi = 0$



Operators splitting

Governing equations

$$\begin{cases} \partial_t \rho + \nabla \cdot (\rho \mathbf{u}) = 0 \\ \partial_t (\rho \mathbf{u}) + \nabla \cdot (\rho \mathbf{u} \otimes \mathbf{u}) + \nabla p = -\rho \nabla \psi + \nabla \cdot \boldsymbol{\tau}_{visc} + \sigma \kappa \mathbf{n} \delta(\phi) \\ \partial_t (\rho E) + \nabla \cdot ((\rho E + p) \mathbf{u}) = -\rho \mathbf{u} \cdot \nabla \psi + \nabla \cdot (\boldsymbol{\tau}_{visc} \mathbf{u}) + \sigma \kappa \mathbf{u} \cdot \mathbf{n} \delta(\phi) + \nabla \cdot (k \nabla T) \\ \partial_t \phi + \mathbf{u} \cdot \nabla \phi = 0 \end{cases}$$

$$E = e + \frac{|\mathbf{u}|^2}{2} \quad \rho e = \frac{p + \gamma \pi^\infty}{\gamma - 1} \quad p = \rho r T - \gamma \pi^\infty \quad r = c_p - c_v \quad \gamma = c_p / c_v$$

ψ : gravitational potential $\boldsymbol{\tau}_{visc}$: viscous tensor

κ : interface Curvature k : thermal conductivity

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Acoustic / Transport / Diffusion decomposition

$$\begin{cases} \partial_t \rho + \rho \nabla \cdot \mathbf{u} + \mathbf{u} \cdot \nabla \rho = 0 \\ \partial_t (\rho \mathbf{u}) + \rho \mathbf{u} \nabla \cdot \mathbf{u} + \nabla p + \mathbf{u} \cdot \nabla (\rho \mathbf{u}) = -\rho \nabla \psi + \sigma \kappa \mathbf{n} \delta(\phi) + \nabla \cdot \boldsymbol{\tau}_{visc} \\ \partial_t (\rho E) + \rho E \nabla \cdot \mathbf{u} + \nabla \cdot (\rho \mathbf{u}) + \mathbf{u} \cdot \nabla (\rho E) = -\rho \mathbf{u} \cdot \nabla \psi + \sigma \kappa \mathbf{u} \cdot \mathbf{n} \delta(\phi) + \nabla \cdot (\boldsymbol{\tau}_{visc} \mathbf{u}) + \nabla \cdot (k \nabla T) \\ \partial_t \phi + \mathbf{u} \cdot \nabla \phi = 0 \end{cases}$$

Numerical resolution

① Acoustic system

- perform a **Suliciu-type relaxation** to deal with any equation of state and obtain **linearly degenerate system** [Bouchut, 2004]
- resolve the Riemann problem that takes the source term (gravity) and jump conditions related to **surface tension & phase change (future work)** into account in a consistent way
- **well-balanced scheme** for gravity and surface tension : discretization at cell face or interface
- **low-Mach correction:** the accuracy of the global scheme independent form Mach number [Chalons et al., 2016]
- **implicit scheme:** $\Delta t = \Delta x / \max(|u|)$ (**future work**)

② Transport system

- sharp interface & avoid mixing cell:ghost fluid method
- Level-Set advection: **One-Step (OS) scheme** [Daru and Tenaud, 2004]
coupled time-space approach: same order accuracy at time & space

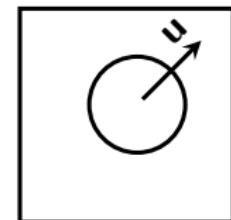
③ Diffusion system (viscous & heat transfer)

- classical finite volume discretization

Numerical test

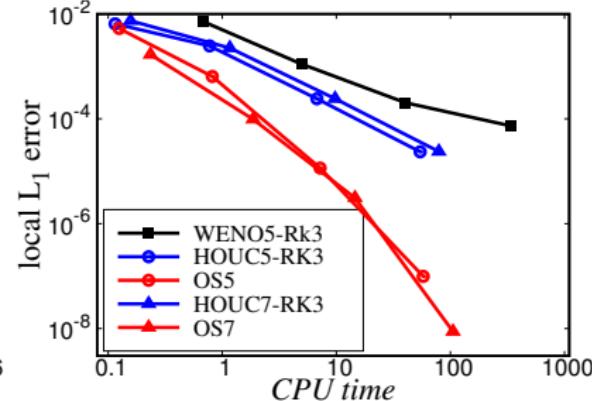
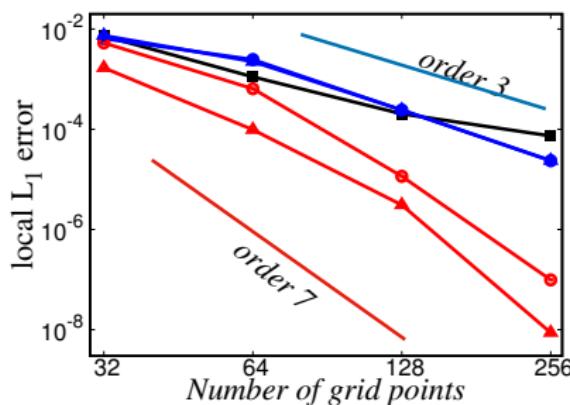
Circle advection

- advection in diagonal direction with a periodic boundary condition
- $t=100$ (100 periods) $CFL=0.5$
- local L_1 error ($|\phi| \leq 3\Delta x$)



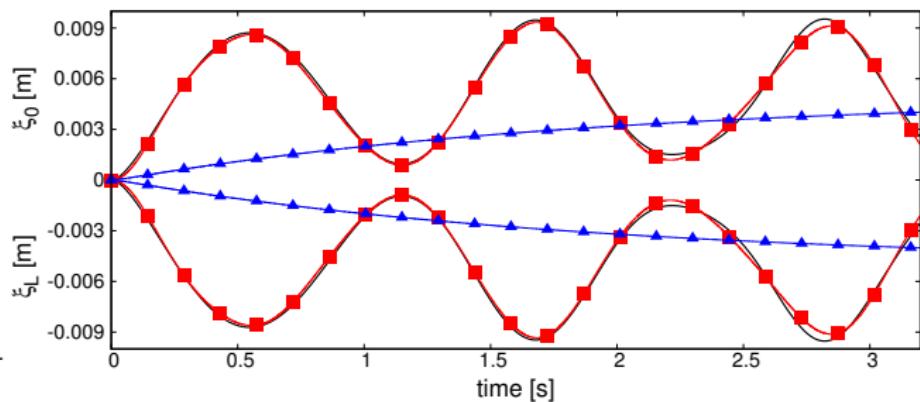
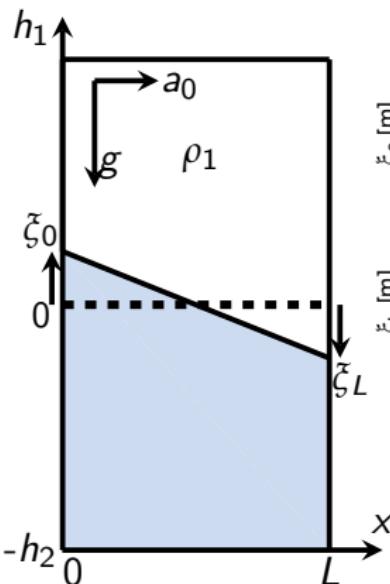
Numerical scheme

- Coupled time-space approach: high-order One-Step (OS) scheme
- Separate time-space approach: WENO-RK [Jiang and Shu, 1996]
HOUC (High-Order upwind centered)-RK [Nourgaliev, 2007]



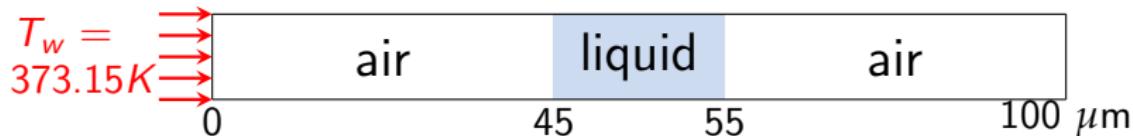
Two-dimensional sloshing [Chanteperrin, 2004]

- $\rho_2/\rho_1 = 1000 \quad h_1/L = 1.25 \quad h_2/L = 1$
- $c_1 \approx 370\text{m/s} \quad c_2 \approx 1500\text{m/s} \quad \text{Ma}_{\max} \approx 2 \times 10^{-5}$
- $a_0/g = 0.01$ resolution : $L/\Delta x = 40$

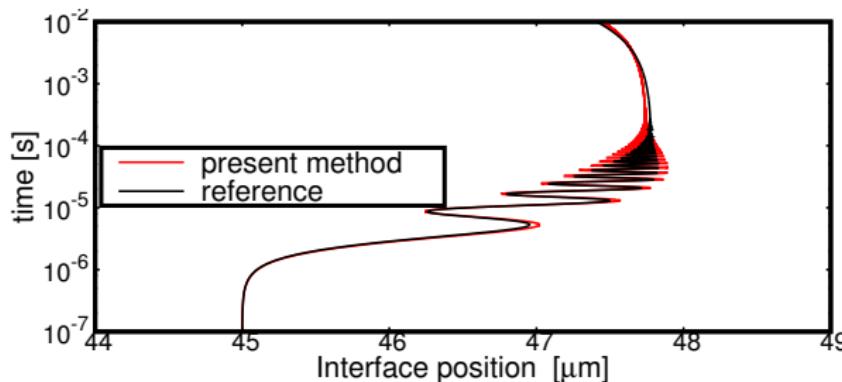


— analytical solution: linearized potential theory
■ with low-Mach correction
▲ without low-Mach correction

1D non-isothermal problem



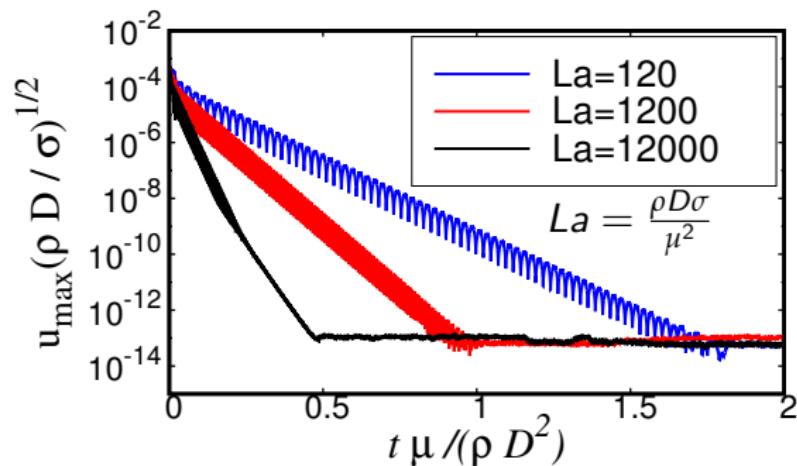
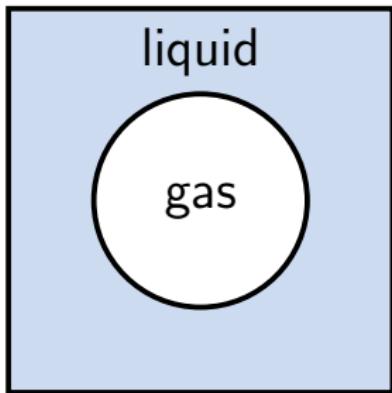
- $T_{liq} = T_{air} = 293.15 \text{ K}$
- liquid: $\rho = 1000 \text{ kg/m}^3$ $c = 1500 \text{ m/s}$ $c_p = 4184 \text{ J K}^{-1} \text{ kg}^{-1}$
 $k = 0.6 \text{ W m}^{-1} \text{ K}^{-1}$ $\mu = 1.82 \times 10^{-5} \text{ Pa s}$
- air: $\rho = 1.204 \text{ kg/m}^3$ $c = 340 \text{ m/s}$ $c_p = 1004.5 \text{ J K}^{-1} \text{ kg}^{-1}$ $k = 0.0256 \text{ W m}^{-1} \text{ K}^{-1}$ $\mu = 0.001 \text{ Pa s}$



Trajectories of left interface: reference [Daru et al., 2010]

Parasitic currents

- liquid $\rho = 1 \quad c = 20 \quad p_l$
- gas $\rho = 1 \quad c = 1 \quad p_g = p_l + \sigma\kappa$

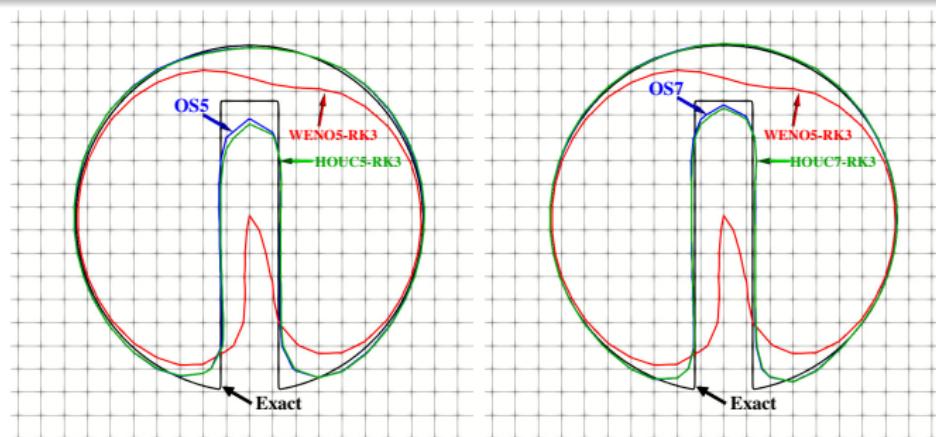


Evolution of velocity fluctuations of a 2D bubble in a slightly compressible liquid with different Laplace numbers by varying the fluid viscosity ($R = 12.8 \triangle x$)

Numerical test

Zalesak disk

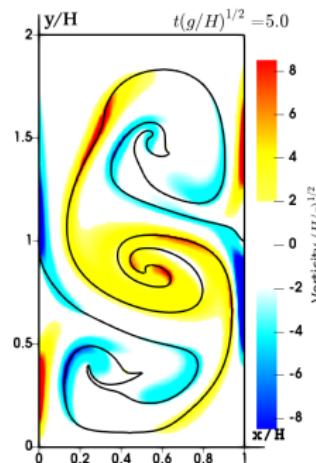
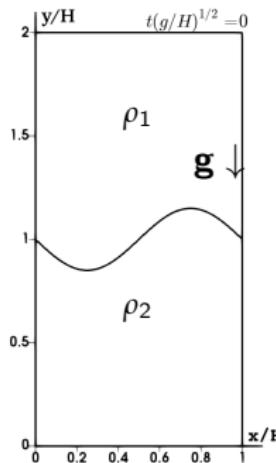
- disk with a rectangular slot
- Rotating velocity field: $u = \pi(50 - y)/314 \quad v = \pi(x - 50)/314$
- one full rotation CFL=0.5



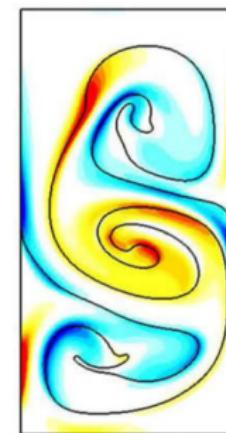
*Interface representation of the Zalesak disk after one full rotation
resolution: 50 × 50*

Rayleigh-Taylor instabilities

- $\rho_1 = 1.8 \quad \rho_2 = 1 \quad Ma_{\max} \approx 10^{-2}$
- $Re = \sqrt{(H/2)^3/g}/\nu = 420 \quad Mesh : 312 \times 624$



(a)



(b)

Vorticity and interface representation: (a) present Level-Set formulation
(b) incompressible Level-Set formulation [Colicchio, 2004]

Thank you for your attention !

References I

-  **Bouchut, F. (2004).**
Springer Science & Business Media.
-  **Chalons, C., Girardin, M., and Kokh, S. (2016).**
An all-regime lagrange-projection like scheme for the gas dynamics equations on unstructured meshes.
Communications in Computational Physics, 20(1):188–233.
-  **Chanteperdrix, G. (2004).**
Modélisation et Simulation Numérique des Ecoulements Diphasiques à Interfaces Libres. Application à l'étude des mouvements de liquides dans les réservoirs de véhicules spatiaux.
PhD thesis, Ecole Nationale Supérieure de l'Aéronautique et de l'Espace.
-  **Colicchio, G. (2004).**
Violent disturbance and fragmentation of free surfaces.
PhD thesis, University of Southampton, UK.
-  **Daru, V. and Tenaud, C. (2004).**
Journal of computational physics.

References II

-  Faccanoni, G., Kokh, S., and Allaire, G. (2012).
Modelling and simulation of liquid-vapor phase transition in compressible flows based on thermodynamical equilibrium*.
ESAIM: Mathematical Modelling and Numerical Analysis, 46(5):1029–1054.
-  Fechter, S., Munz, C.-D., Rohde, C., and Zeiler, C. (2017).
A sharp interface method for compressible liquid–vapor flow with phase transition and surface tension.
Journal of Computational Physics, 336:347–374.
-  Jiang, G.-S. and Shu, C.-W. (1996).
Journal of computational physics.
-  Nourgaliev, R. R. (2007).
Journal of Computational Physics.
-  Osher, S. and Sethian, J. A. (1988).
Fronts propagating with curvature-dependent speed: algorithms based on hamilton-jacobi formulations.
Journal of computational physics, 79(1):12–49.

References III

-  Sato, Y. and Ničeno, B. (2013).
A sharp-interface phase change model for a mass-conservative interface tracking method.
Journal of Computational Physics, 249:127–161.
-  Tanguy, S., Sagan, M., Lalanne, B., Couderc, F., and Colin, C. (2014).
Benchmarks and numerical methods for the simulation of boiling flows.
Journal of Computational Physics, 264:1–22.